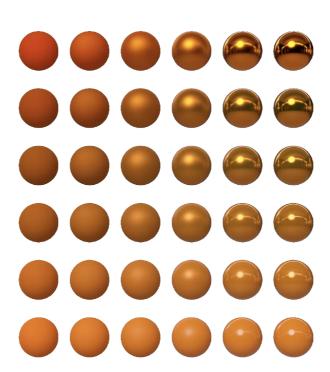


# Advanced Computer Graphics Physically-Based Lighting/Rendering



G. Zachmann
University of Bremen, Germany
cgvr.cs.uni-bremen.de





# Radiometry: Physically Measuring Light



- Interaction of light with objects much larger than its wavelength
- Geometric optics suffices for most cases in computer graphics
  - No interference, no dispersion, etc.
- Advantages:
  - Linearity: effect(light 1 + light 2) = effect(light 1) + effect(light 2)
  - Energy conservation
  - No polarization
  - Lights of different wavelengths are independent of each other (no fluorescence)

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# Four Radiometric Quantities



- Flux: how much energy flows from/to/through a surface per unit of time
  - Think of it as "photons per second"
  - Symbol:  $\Phi$ , units = Watt
  - Example light bulb: surface = sphere around light bulb
    - Radius doesn't matter: #photons crossing the boundary of the sphere is the same
  - Example table surface: doesn't matter where photons come from





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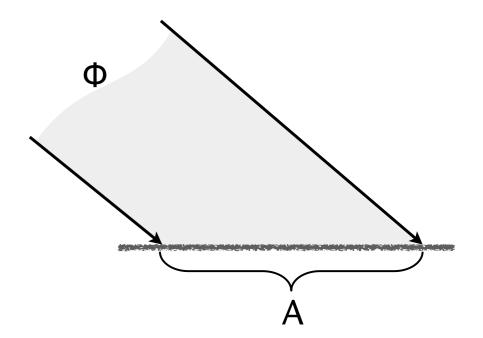
Irradiance: energy density = flux per area

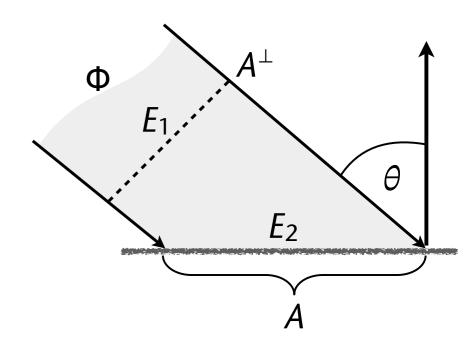
$$E = \frac{d\Phi}{dA}$$

- Example: 50 Watts incident on table surface of 2m<sup>2</sup> -> each point on the table surface receives irradiance of 25 Watt/m<sup>2</sup>
- Lambert's law: irradiances on surfaces are related by the  $cos(\theta)$  factor

$$E_1 = rac{\mathrm{d} \Phi}{\mathrm{d} A^{\perp}}$$
  $E_2 = rac{\mathrm{d} \Phi}{\mathrm{d} A} = rac{\mathrm{d} \Phi \cos heta}{\mathrm{d} A^{\perp}} = E_1 \cos heta$ 

since 
$$A^{\perp} = A \cos \theta$$



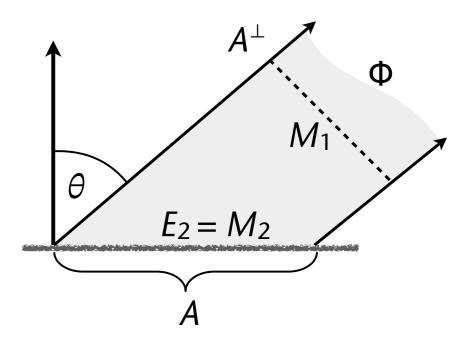






• Exitance (a.k.a. exiting irradiance, radiosity): same thing as irradiance, except for light leaving a surface

$$M_1 = \frac{\mathrm{d}\Phi}{\mathrm{d}A^{\perp}} = \frac{\mathrm{d}\Phi}{\cos\theta\mathrm{d}A} = \frac{E_2}{\cos\theta}$$

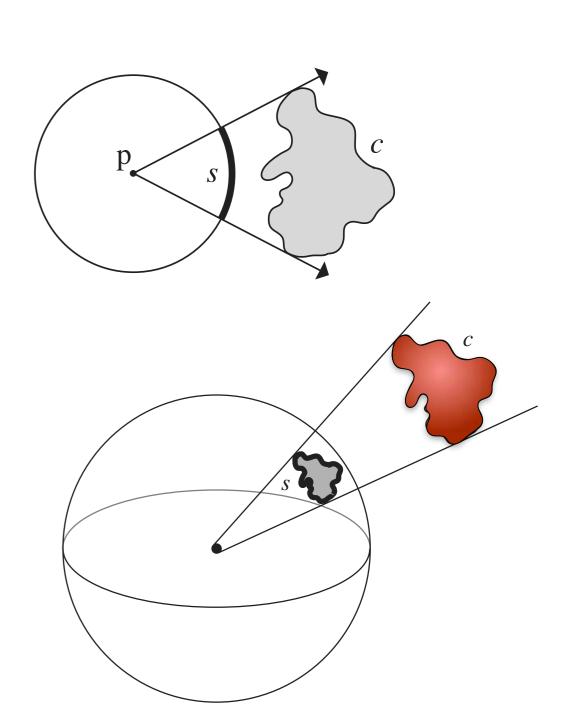




# Solid Angles



- Generalization of angles to 3D:
  - Analog: subtended angle of an object in 2D = length of arc of its projection on the unit circle
    - Units = radians (rad), range =  $[0,2\pi]$
  - Subtended solid angle of an object in 3D = area of its projection on the unit sphere
    - Units = steradians (sr), range =  $[0,4\pi]$
    - Solid angle = area on the unit sphere
- Notation:
  - $^{\prime}\cosarphi$  sin  $heta^{ee}$ •  $\omega = \text{direction} = \text{unit vector} = \begin{pmatrix} \sin \varphi \sin \theta \\ \cos \theta \end{pmatrix}$
  - $d\omega$  = infinitesimal solid angle in direction  $\omega$
  - Sometimes,  $\omega$  is used to denote a solid angle, too



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#### Note on the Side: How to Make a Change of Variables in Integrations

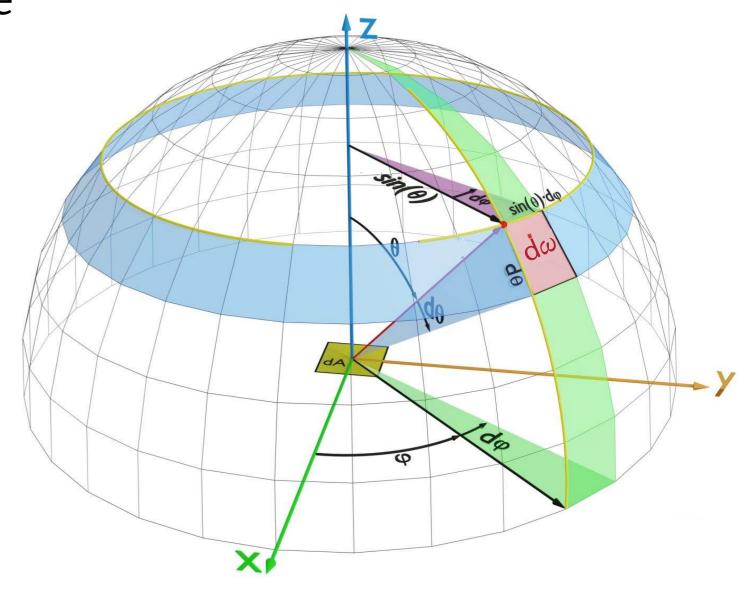


- When evaluating integrals over the hemisphere,  $\Omega$ , we could use the solid angle directly, or polar angles instead
  - Polar angles:  $\varphi \in [0, 2\pi]$  ,  $\theta \in [0, \pi/2]$  where  $\theta = 0$  points *upwards*
  - Consider a square solid angle,  $d\omega$ :

$$d\omega = width \cdot height$$
$$= (\sin \theta \cdot d\varphi) \cdot d\theta$$

Therefore,

$$\int_{\Omega} \dots d\omega = \int_{0}^{\pi/2} \int_{0}^{2\pi} \dots \sin(\theta) d\varphi d\theta$$



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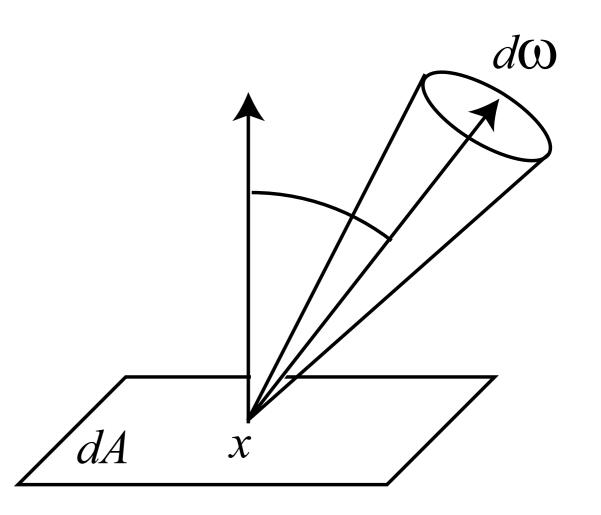




Intensity: flux over solid angles

$$I=rac{\mathsf{d}oldsymbol{\phi}}{\mathsf{d}\omega}$$

- Think "number of photons sent into a small solid angle around direction  $\omega$ "
- Only meaningful for a outgoing light emanating from a single pont x, or light coming in from d $\omega$  and converging in point x
- Like irradiance, it is a density, but a different kind of energy density!



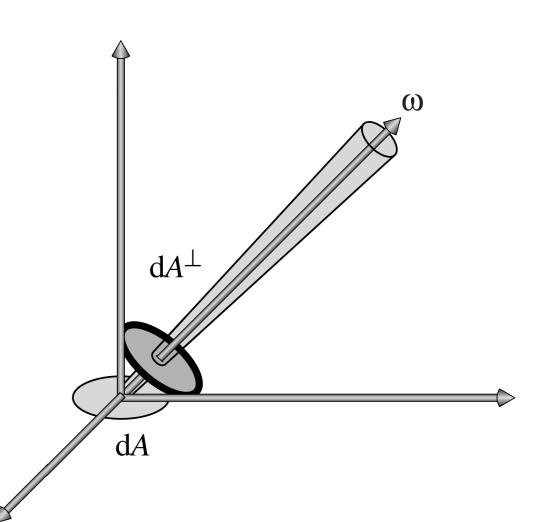




Radiance: "how much power arrives at (or leaves from)
 a specific point on a surface, per unit solid angle, and
 per unit projected area"

$$L = \frac{\mathrm{d}^2 \Phi}{\mathrm{d}\omega \, \mathrm{d}A^{\perp}}$$

- Think: "limit of incident light at the surface as a cone of incident directions, d $\omega$ , becomes very small, and as the local area on the surface, dA, also becomes very small"
- Think: combination of 2 densities, irradiance and intensity
- Think: a neuron on the retina, or a pixel on a CCD chip inside a camera, measures radiance (at a specific small area, arriving from a specific small cone)



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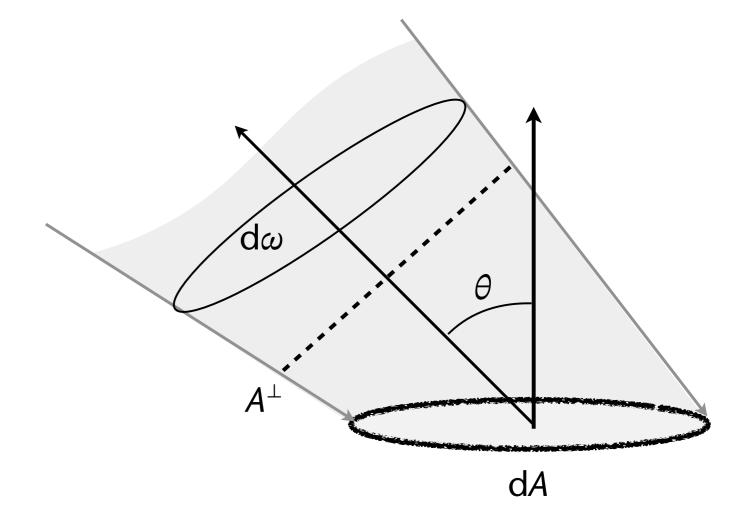
## Simple Properties and Relationships of Radiance



• Relationships between *L*, *E*, *I*:

$$L = \frac{\mathrm{d}^2 \Phi}{\mathrm{d}\omega \, \mathrm{d}A^{\perp}} = \frac{\mathrm{d}E}{\mathrm{d}\omega} = \frac{\mathrm{d}I}{\mathrm{d}A^{\perp}}$$

$$d^2 \Phi = \frac{d^2 \Phi}{\cos \theta \, dA \, d\omega}$$



- Notes:
  - L is a 5-dimensional function  $L(p, \omega)$
  - Radiance is defined both for incoming,  $L_i(p,\omega)$ , as well as outgoing light,  $L_o(p,\omega)$ !
  - Direction  $\omega$  always points away from point p
  - For incoming light ( $L_i$ ):  $\cos \theta = \mathbf{n} \cdot \mathbf{l}$ , for outgoing light ( $L_o$ ):  $\cos \theta = \mathbf{n} \cdot \mathbf{v}$



# Important Property: Radiance is Invariant Along Straight Paths



- Arguably the most important property
- Consider light that travels from point p to point q (without any participating media)
- Exiting radiance

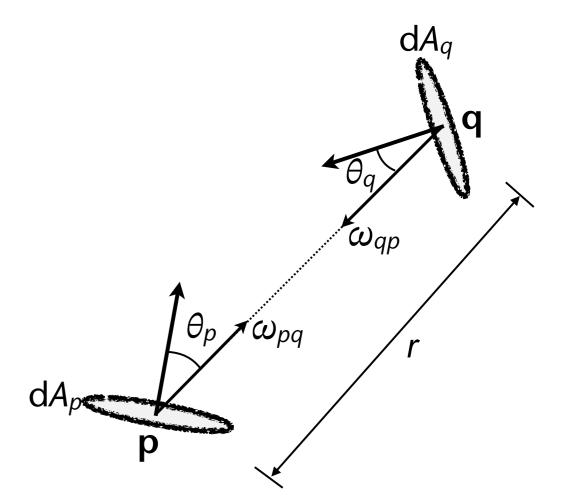
$$L_o(\omega_{pq}) = \frac{d^2 \phi_{pq}}{\cos \theta_p dA_p d\omega_{pq}}$$
 (1)



$$L_i(\omega_{qp}) = \frac{d^2 \Phi_{qp}}{\cos \theta_q \, dA_q \, d\omega_{qp}}$$
 (2)



$$d^2 \Phi_{pq} = d^2 \Phi_{qp} \tag{3}$$







• Solve (1) and (2) for  $d^2\phi$  and plug into (3):

$$L_o(\omega_{pq})\cos\theta_p\,\mathrm{d}A_p\,\mathrm{d}\omega_{pq} = L_i(\omega_{qp})\cos\theta_q\,\mathrm{d}A_q\,\mathrm{d}\omega_{qp} \tag{4}$$

- Compute  $d\omega_{pq}=rac{dA^\perp}{r^2}=rac{\cos\theta_q\,dA_q}{r^2}$  , and similarly  $d\omega_{qp}$  , and plug into (4)
- Thus,

$$L_o(\omega_{pq})\cos\theta_p\,\mathrm{d}A_p\,\frac{\cos\theta_q\,\mathrm{d}A_q}{r^2}=L_i(\omega_{qp})\cos\theta_q\,\mathrm{d}A_q\,\frac{\cos\theta_p\,\mathrm{d}A_p}{r^2}$$

And, hence,

$$L_o(\omega_{pq}) = L_i(\omega_{qp})$$

• Thus, radiances are the physically correct quantities to exchange between points!



### Reasons for Radiance



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- Eyes / cameras measure radiance
- Radiance is invariant w.r.t. distance
- Therefore, physically-correct lighting must compute radiances



### Determining Useful Quantities from Radiance



- Computing the total irradiance at a point on a surface:
  - Irradiance from a specific direction:

$$L_i(\omega) = \frac{d^2 \Phi(\omega)}{\cos \theta \, dA \, d\omega} \qquad L_i(\omega) \cos \theta \, d\omega = \frac{d \Phi(\omega)}{dA} \quad = E(\omega)$$

• Needs to be integrated over all directions to obtain the total "brightness" (Watt/m²):

$$E = \frac{\mathrm{d}\Phi}{\mathrm{d}A} = \int_{\Omega} L_i(\omega) \cos\theta \,\mathrm{d}\omega$$

where  $\Omega$  is the hemisphere centered on the normal in point p





- Irradiance at a point p arriving from a distant surface A:
  - From the definition of radiance:

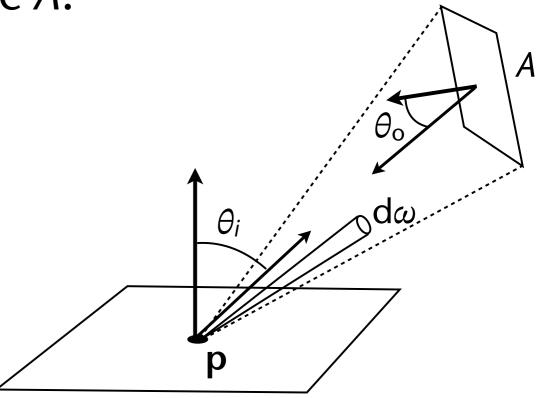
$$dE = L \cos \theta_i d\omega$$

Integrating leads to:

$$E = \int_{\Omega'} L(\theta_i, \varphi_i) \cos \theta_i \, \mathrm{d}\omega$$

- However, this would need to be integrated over that part of  $\Omega$  where A is visible  $\to \Omega'$ , which is cumbersome
- Trick: replace  $d\omega$  by  $d\omega = \cos\theta_o \cdot \frac{dA}{r^2}$
- Thus,

$$E = \int_{A} L(\omega_{i}) \cos \theta_{i} \cos \theta_{o} \frac{1}{r^{2}} dA$$







- Total flux (power) emitted from a surface:
  - Consider the outgoing radiance from a surface A

$$L_o(\omega) = \frac{d^2 \Phi(\omega)}{\cos \theta \, dA \, d\omega} \qquad L_o(\omega) \cos \theta \, dA \, d\omega = d^2 \Phi(\omega)$$

- Note that  $\theta$  is one of the polar angles of  $\omega$ , and  $\cos \theta = \mathbf{n} \cdot \mathbf{v}$ , where  $\mathbf{v}$  points in the direction of  $\omega$
- Integrating over all directions and the whole surface yields

$$\Phi_o = \int_A \int_C L_o(\omega) \cos \theta \, d\omega \, dA$$

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• Except: we don't know  $L_0$  ... (yet)



### Example



- Consider area light source  $A = 10x10 \text{ cm}^2$
- Assume, all points on A have radiance  $L_o(\phi, \theta) = 6000 \cos \theta \frac{W}{cr_b m^2}$
- Exiting power of the light source:

$$\Phi = \int_{A} \int_{\Omega} L_{o} \cos \theta \, d\omega \, dA = \int_{A} \int_{\Omega} 6000 \cos^{2} \theta \, d\omega \, dA$$

$$= \int_{A} 6000 \int_{0}^{2\pi} \int_{0}^{\pi/2} \cos^{2}\theta \sin\theta \,d\theta \,d\varphi \,dA = \int_{A} 6000 \cdot 2\pi \cdot \left[ \frac{-\cos^{3}\theta}{3} \right]_{0}^{\pi/2} \,dA$$

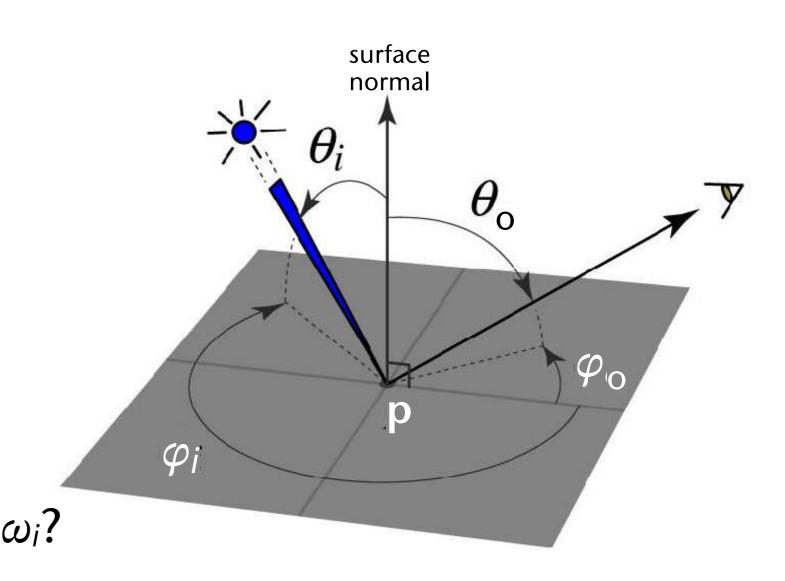
$$=6000 \cdot 2\pi \cdot \frac{1}{3} \frac{W}{m^2} \cdot 0.1 m \cdot 0.1 m \approx 125 W$$



# The BRDF: Interaction of Light with Surfaces



- Given some light shining on point **p** on a surface, coming in from direction  $\omega_i$ , how much light exits that point into another direction  $\omega_o$ ?
- Assumptions:
  - Light does not change wavelength
  - Reflection / scattering is instantaneous
  - Light does not travel inside the material (no subsurface scattering)
- Question: how much radiance is exiting in direction  $\omega_0$ , as a result of incident radiance coming from direction  $\omega_i$ ?





#### Definition of BRDF



• Irradiance at point p depends on incident radiance,  $L_i$ :

$$dE(\omega_i) = L_i(\omega_i) \cos \theta_i d\omega_i$$

- Some of that irradiance gets scattered into direction  $\omega_o$ , leading to exiting radiance  $L_o$
- By the linearity assumption:  $dL_o(\omega_o) \propto dE(\omega_i)$
- The proportionality factor is defined as the Bidirectional Reflectance Distribution Function (BRDF)

$$\rho(\omega_o, \omega_i) = \frac{dL_o(\omega_o)}{dE(\omega_i)} = \frac{dL_o(\omega_o)}{L_i(\omega_i)\cos\theta_i d\omega_i}$$



#### Note on Notations of the BRDF Function



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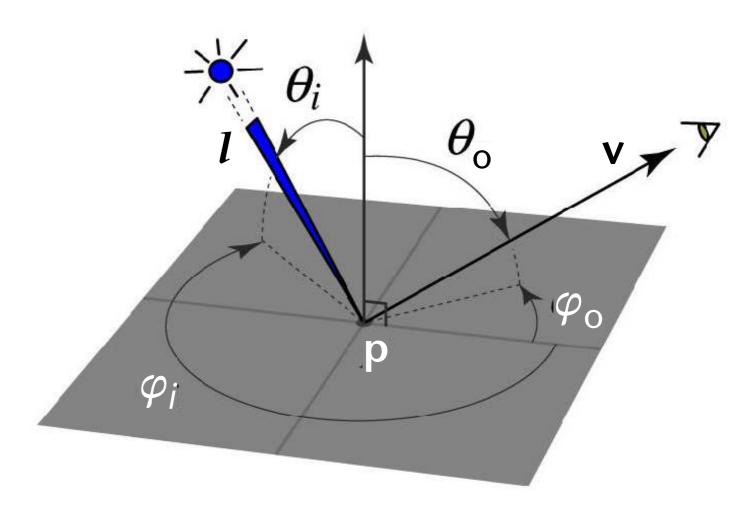
The BRDF function is a 4D function:

$$\rho(\omega_o, \omega_i) = \rho(\theta_o, \varphi_o, \theta_i, \varphi_i)$$

- Assumes that  $\rho$  is constant over the surface
- Sometimes,  $\rho$  is written in terms of vectors:

$$ho=
ho(m{l},m{v})$$

In practice, this version is preferred,
 since cos(angle) = scalar product of vectors





# Properties of the BRDF



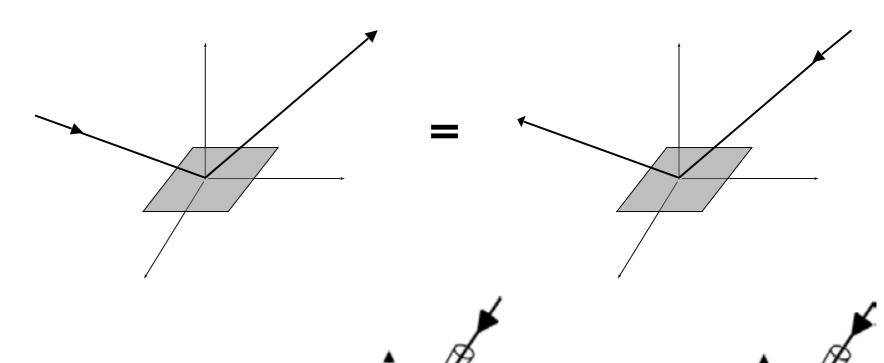
• Reciprocity:

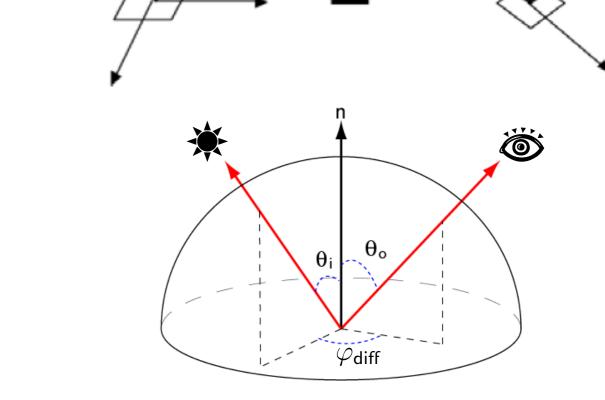
$$\rho(\omega_i, \omega_o) = \rho(\omega_o, \omega_i)$$



- Examples: brushed metal, hair, satin, ...
- For many materials, though, it is isotropic
  - Consequence,  $\rho = \rho(\theta_i, \theta_o, \varphi_{\text{diff}})$
- Positivity:  $\rho \ge 0$  everywhere

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# The Reflectance Equation (aka Rendering Equation)



- Computing the outgoing radiance
- Solve definition of BRDF for  $L_0$ :

$$dL_o(\omega_o) = \rho(\omega_o, \omega_i) L_i(\omega_i) \cos \theta_i d\omega_i$$

Integrate over all incoming directions:

$$L_o(\omega_o) = \int_{\Omega} \rho(\omega_o, \omega_i) L_i(\omega_i) \cos \theta_i \, d\omega_i$$

- In practice, we sample into "all" possible directions  $\omega_i$ :
  - Many different, clever methods to choose the "best" sampling directions (Monte) Carlo, variance reduction, importance-based sampling)





• Alternative "spelling" of the reflectance equation:

$$L_o(\mathbf{v}) = \int_{\Omega} \rho(\mathbf{l}, \mathbf{v}) \cdot L_i(\mathbf{l}) (\mathbf{n} \cdot \mathbf{l}) \, \mathrm{d}\omega_i$$

• Note that the term  $\cos \theta_i = \mathbf{n} \cdot \mathbf{l}$  is *not* part of the BRDF!

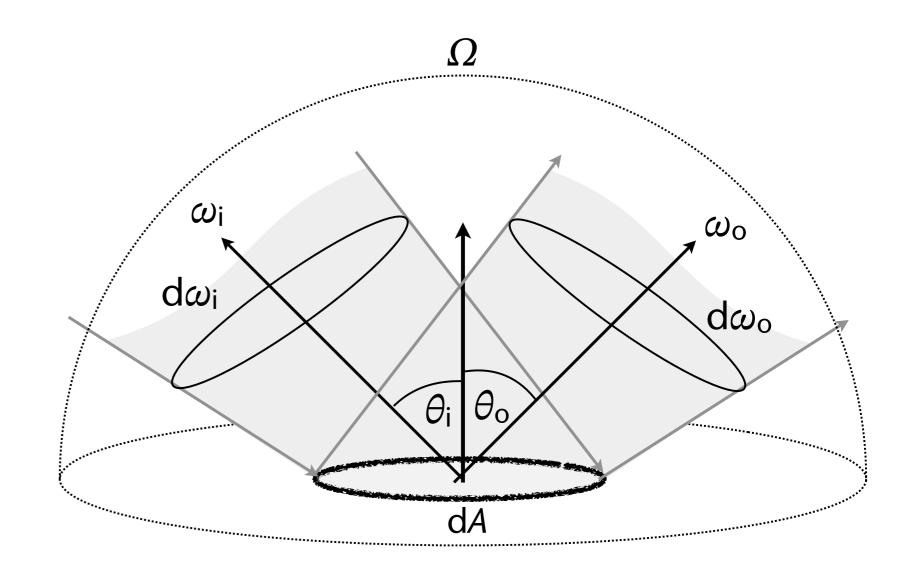


# Another Property of BRDFs: Conservation of Energy



- No energy can be created
  - (Also, it cannot be destroyed, but that is usually modeled differently in CG)
- Incoming (incident) irradiance from all directions:

$$E = \int_{\Omega} L_i(\omega_i) \cos \theta_i \, \mathrm{d}\omega_i$$



Outgoing (exitant) irradiance in all dir's:

$$M = \int_{\Omega} L_o(\omega_o) \cos \theta_o \, d\omega_o = \int_{\Omega} \int_{\Omega} \rho(\omega_o, \omega_i) L_i(\omega_i) \cos \theta_i \, d\omega_i \, \cos \theta_o \, d\omega_o$$





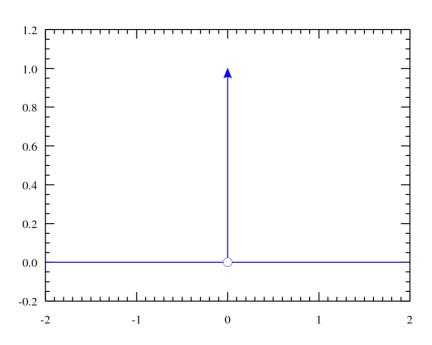
- Because of conservation of energy,  $M \stackrel{!}{\leq} E$ , i.e.,  $\frac{M}{E} \stackrel{!}{\leq} 1$  for all and  $any L_i$  (!)
- Note:  $\frac{M}{E}$  is sometimes called "reflectance"
  - Meaning: how much of all incoming light (from all directions) is reflected (scattered) into all directions
- Choose the special incoming radiance  $L_i(\omega_i) = \bar{L}_i \delta(\omega_i \bar{\omega})$  where  $\delta$  is the Dirac function

$$\delta(x) = \begin{cases} +\infty & x = 0 \\ 0 & x \neq 0 \end{cases}$$

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Important property:

$$\int f(x)\delta(x-\bar{x})\,\mathrm{d}x=f(\bar{x})$$







• With this special  $L_i$ , the incoming irradiance becomes

$$E = \int_{\mathcal{O}} \bar{L}_i \delta(\omega_i - \bar{\omega}) \cos \theta_i \, d\omega_i = \bar{L}_i \cos \bar{\theta}$$
 (2)

where  $\bar{L}_i$  is coming in from only one single direction  $\bar{\omega} = (\bar{\theta}, \bar{\phi})$ 

The outgoing irradiance becomes

$$M = \int_{\Omega} \int_{\Omega} \rho(\omega_o, \omega_i) \bar{L}_i \delta(\omega_i - \bar{\omega}) \cos \theta_i \cos \theta_o \, d\omega_i \, d\omega_o$$

$$= \int_{O} \rho(\omega_{o}, \bar{\omega}) \bar{L}_{i} \cos \bar{\theta} \cos \theta_{o} d\omega_{o}$$
 (3)





Plugging (2) and (3) into (1) gives:

$$\frac{M}{E} = \frac{\int_{\Omega} \rho(\omega_o, \bar{\omega}) \bar{L}_i \cos \bar{\theta} \cos \theta_o \, d\omega_o}{\bar{L}_i \cos \bar{\theta}} = \int_{\Omega} \rho(\omega_o, \bar{\omega}) \cos \theta_o \, d\omega_o \leq 1$$

- Remember that energy conservation holds for each and every  $L_i$
- In total, we have the following necessary condition for BRDF's:

$$\forall \omega_i : \int_{\Omega} \rho(\omega_o, \omega_i) \cos \theta_o \, \mathrm{d}\omega_o \stackrel{!}{\leq} 1$$

• This will be important to check for concrete functions  $\rho$ !

Physically-Based Lighting

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### Practical Computation of $L_0$ in Case of Illumination with a Point Light

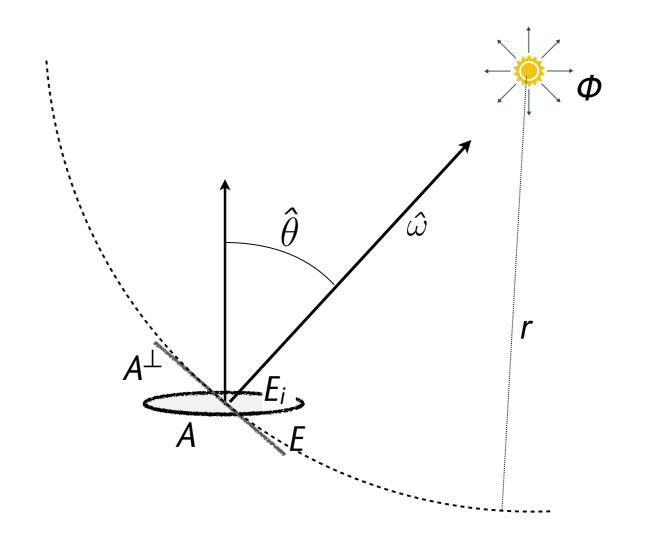


- Note: in real life, point lights never occur (except for stars)
- Assumption: the point light shines in all directions isotropically, with flux  $\phi$
- First: compute irradiance at A
  - Irradiance at  $A^{\perp}$ :

$$E = \frac{\Phi_i}{A^{\perp}} = \frac{\Phi}{4\pi r^2}$$

• Irradiance on A (see slide 4):

$$E(\hat{\omega}) = \cos \hat{\theta} \cdot E = \cos \hat{\theta} \cdot \frac{\Phi}{4\pi r^2}$$







Start with definition of radiance (slide 10):

$$L_i = \frac{\mathrm{d}^2 \Phi}{\cos \theta_i \, \mathrm{d} A \, \mathrm{d} \omega_i}$$

• Resolve to get dE (slide 4):

$$\Rightarrow L_i \cos \theta_i \, d\omega_i = \frac{d^2 \Phi}{dA} = dE$$

Use reflectance equation and insert previous eq.:

$$L_o = \int_{\Omega} \rho L_i \cos \theta_i \, d\omega_i = \int_{\Omega} \rho \, dE(\omega_i)$$





- Since it's a point light source, we receive light from exactly one direction  $\hat{\omega}$ 
  - So, use the "trick" with the Dirac delta function again
- Outgoing radiance can be written as:

$$L_o = \int_{\Omega} \rho \delta(\omega_i - \hat{\omega}) \, dE(\omega_i) = \rho E(\hat{\omega}) = \rho \frac{\cos \hat{\theta}}{4\pi r^2} \Phi$$
(slide 28)

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# Constructing Actual BRDF's



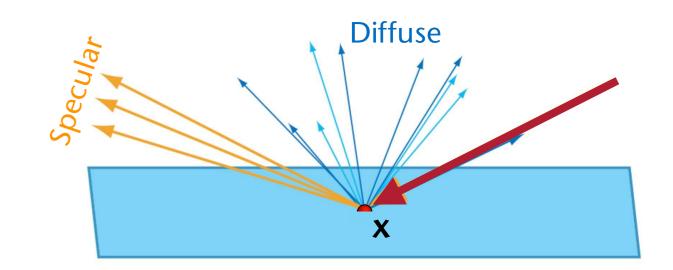
- Two types of reflections: diffuse and specular
- Accordingly, a BRDF consist of

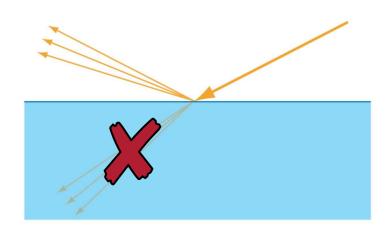
$$\rho = \rho_{\mathsf{diff}} + \rho_{\mathsf{spec}}$$

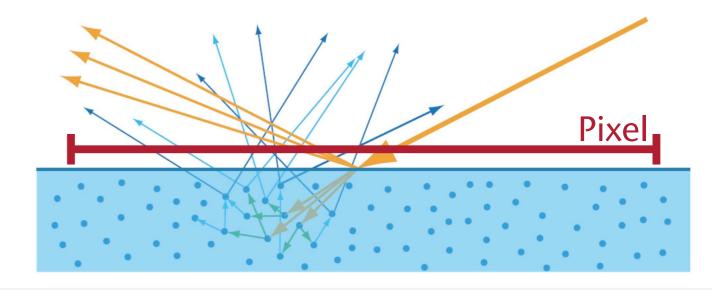
The diffuse term is simply the Lambert model:

$$ho_{\mathsf{diff}}(oldsymbol{l}, \mathbf{v}) = rac{1}{\pi} C_{\mathsf{base}}$$

- Two types of materials: metallic and non-metallic (aka. dielectric)
  - Metallic: no diffuse term
  - Non-metallic: diffuse term due to internal scattering, but travel distance « pixel size





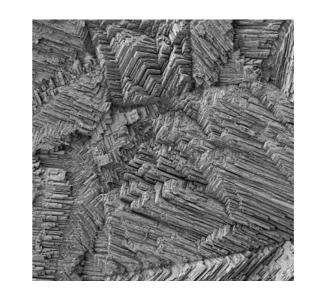


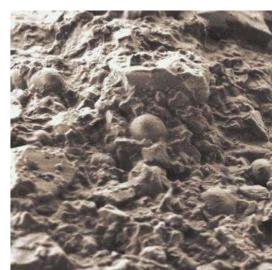


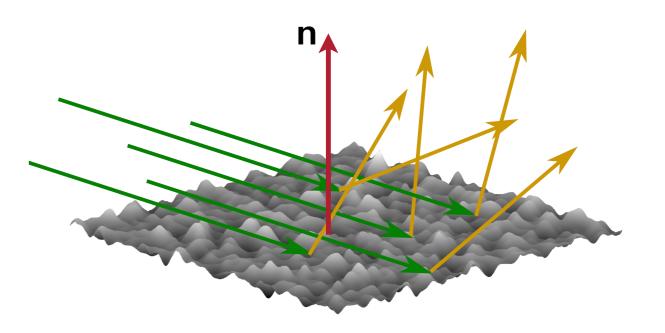
# Constructing the Specular Term



- Usually based on the micro-facet theory
- Even (apparently) smooth surfaces consist of microscopic facets
- Each micro-facet acts like a small perfect mirror
- Macroscopic surface normal n = "average" of all micro-facet normals, m
- Variance of all micro-facets around n depends on (or determines) the roughness of the surface





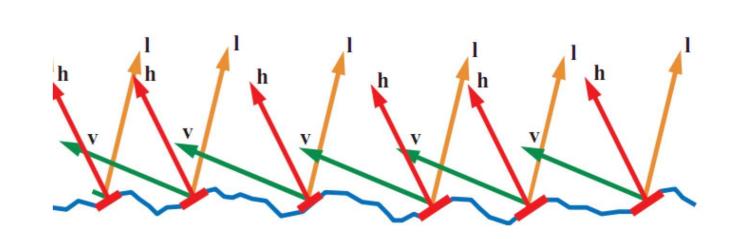


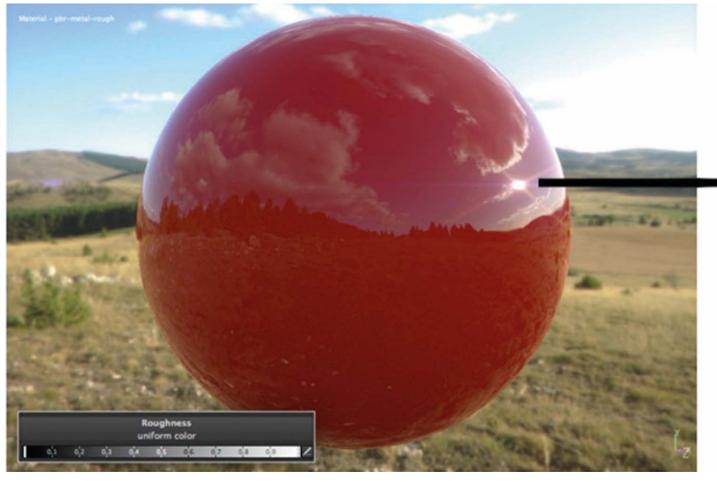


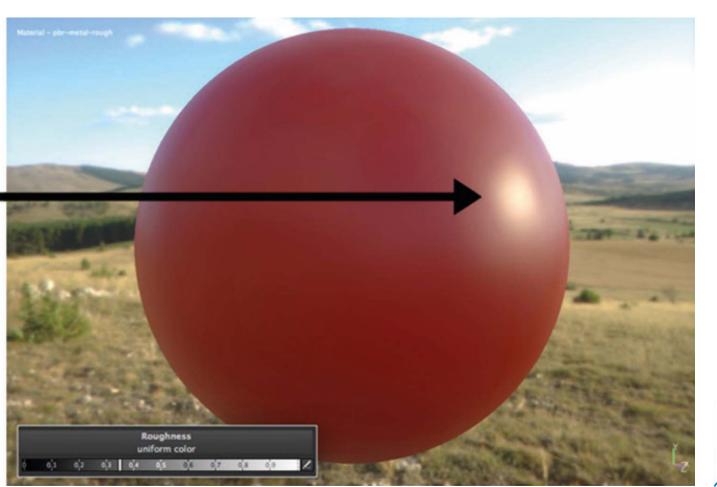
#### 1. The Normal Distribution Function

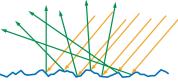


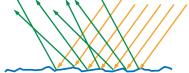
- For given  $l(\omega_i)$  and  $\mathbf{v}(\omega_o)$ , we are only interested in micro-facets with  $\mathbf{m} = \mathbf{h}$
- → Normal distribution function (NDF), D(h), describes "roughness"









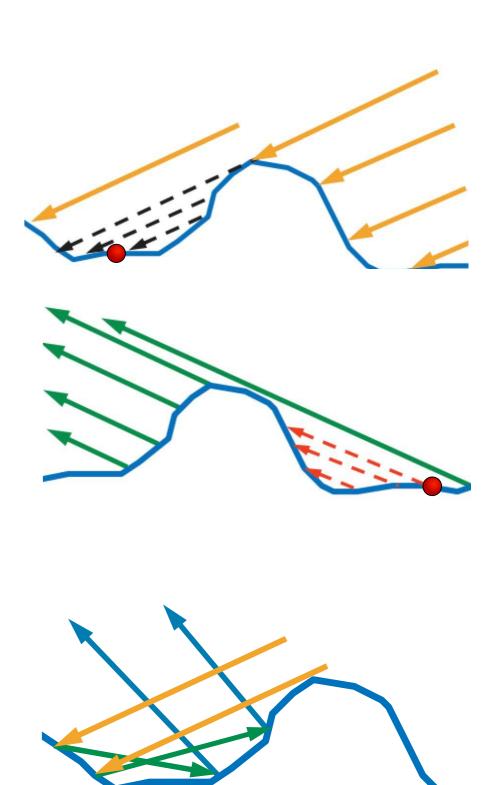




## 2. Geometry Function



- Shadowing and masking effects:
  - Other micro-facets, i.e. with  $\mathbf{m} \neq \mathbf{h}$ , can block either incoming light (shadowing) or are invisible from the outgoing direction (masking)
- Model these effects by a geometry function G(l, v, h)
  - Tells the percentage of surface points with m = h that are not shadowed or masked, as a function of the light direction *l* and the view direction **v**
  - Inter-reflections are (usually) ignored

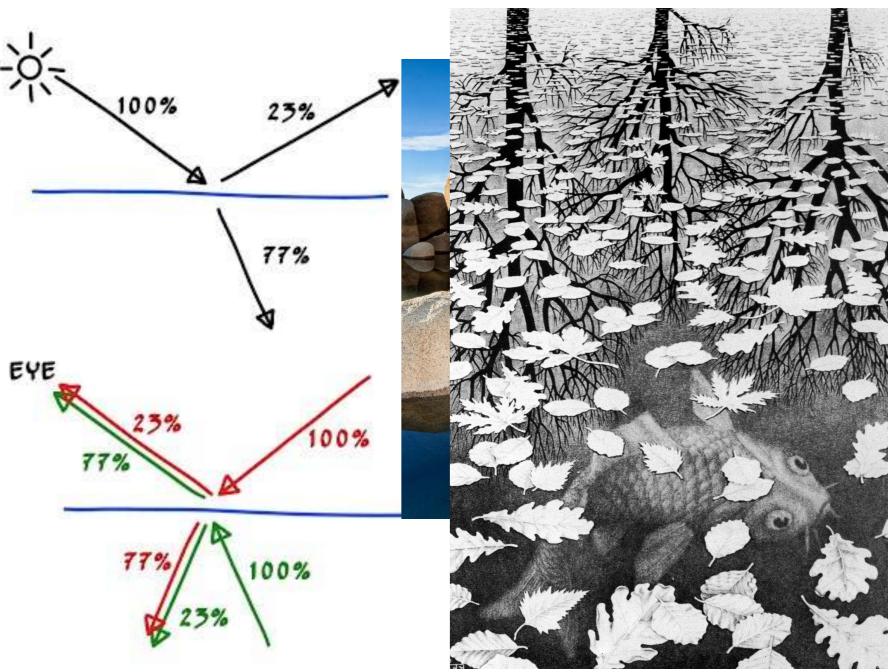




#### 3. Fresnel Function



- Remember the Fresnel effect:
  - Some percentage of light (*F*) gets reflected, the rest (1-*F*) penetrates the surfaces
  - For surfaces between perfectly transparent materials, both incoming lights add nicely up to 100%
  - For opaque materials: the "rest" (1-F) gets absorbed
    - For semi-transparent materials: inbetween; probably best left to the artist
- Modeled by the Fresnel function  $F(\theta)$



[M. C. Escher]



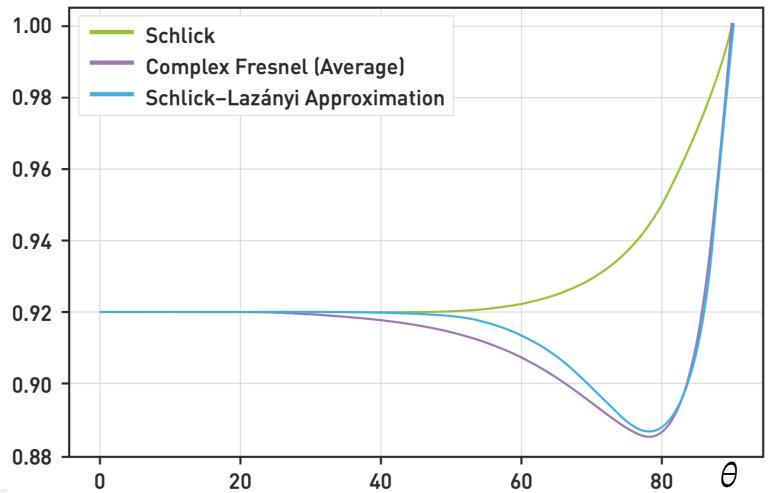
### The Schlick-Lazányi Approximation For Metals



For metals, a more complex approximation should be used:

$$\rho(\theta) \approx \rho_0 + (1 - \rho_0) (1 - \cos \theta)^5 - a \cos \theta (1 - \cos \theta)^6$$

- $\rho$  = fraction of reflected light, rest of light gets absorbed
- Also,  $\rho_0$  depends on the wavelength! (not so for dielectrics)
- Comparison:



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Aluminum, incident light at 450 nm

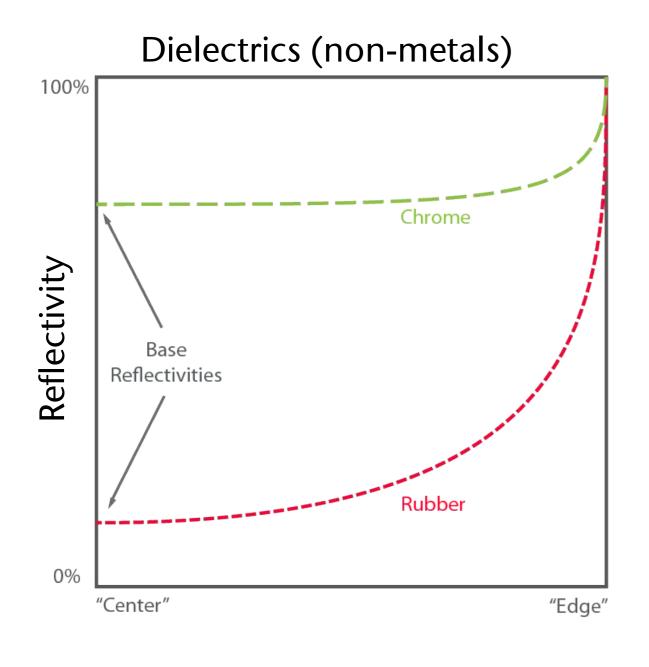


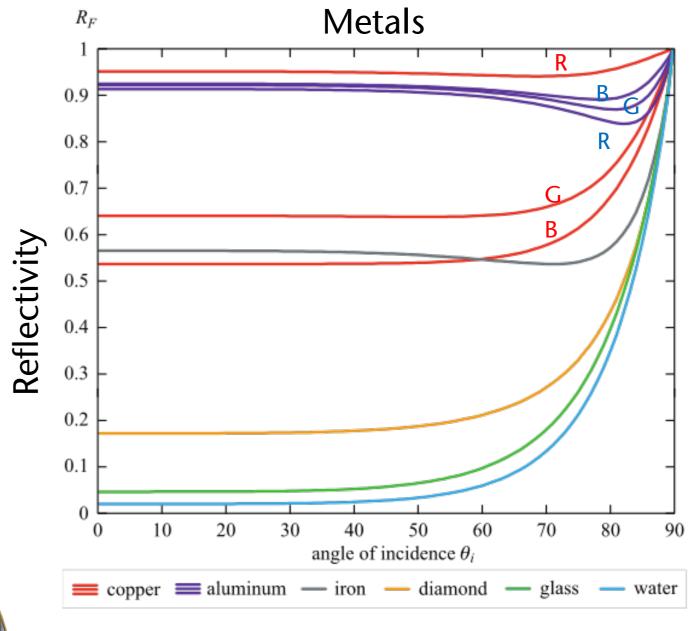


#### Difference between Dielectrics and Metals



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Notice: different curves for different wavelengths!



## The Micro-Facet BRDF as a Whole



Just multiply all three terms (and normalize, so the conditions are met)

$$\rho_{\text{spec}}(\boldsymbol{l}, \mathbf{v}) = \frac{F(\boldsymbol{l}, \mathbf{h})D(\mathbf{h})G(\boldsymbol{l}, \mathbf{v})}{4(\mathbf{n} \cdot \boldsymbol{l})(\mathbf{n} \cdot \mathbf{v})}$$

- Note: evaluate only, if  $\mathbf{n} \cdot \mathbf{l} \geq 0$  and  $\mathbf{n} \cdot \mathbf{v} \geq 0$  (otherwise, either light source, or outgoing direction is "below" surface)
- Observe: when denominator approaches 0, so does G
  - Just check the definition of *G*
- Note: there are a lot of variants, this is just the bare concept



## Approximation of the Fresnel Equation

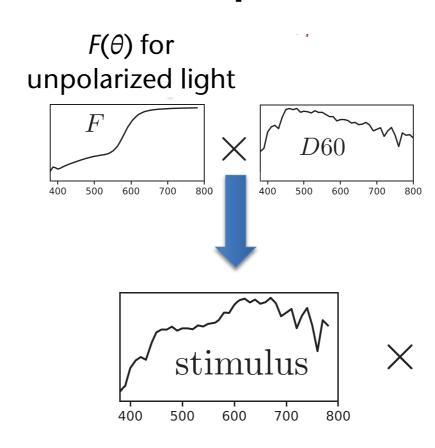


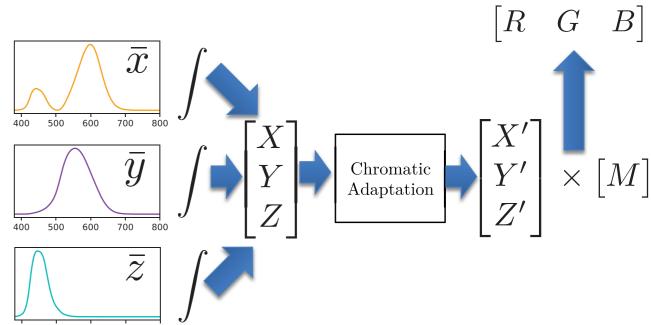
- Physically correct model is extremely complex
  - Depends on wavelength
  - Depends on polarization
- Good approximation is the one by Schlick:

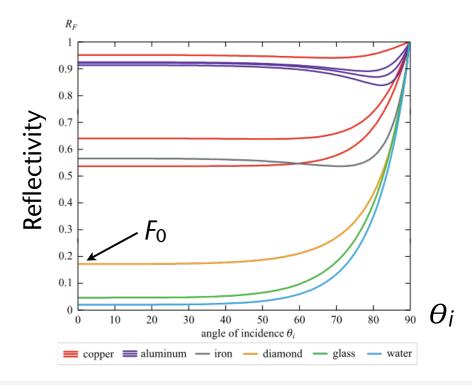
$$F(l, \mathbf{n}) = F_0 + (1 - F_0)(1 - l \cdot \mathbf{n})^5$$

- In micro-facet models: only micro-facets with  $\mathbf{m} = \mathbf{n}$  and  $\mathbf{m} = \mathbf{h} = \frac{1}{2}(\mathbf{l} + \mathbf{v})$
- So, in the micro-facet BRDF:

$$F(l, h) = F_0 + (1 - F_0)(1 - l \cdot h)^5$$









#### Visualization of the Fresnel Term





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Physically-Based Lighting



# The Normal Distribution Function in Detail

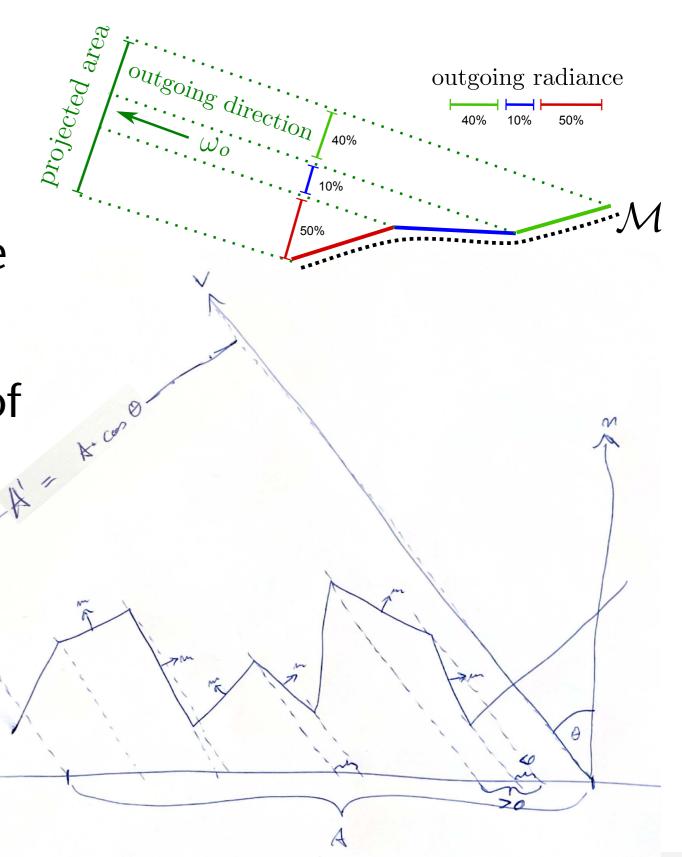


• Roughness is the parameter in any NDF, *D*, determining the variance of the micro-normals around the macroscopic normal

• Intuitive interpretation:  $D(\mathbf{h}) = \text{percentage of the}$ micro-surface  $\mathcal{M}$  that has micro-normals  $\mathbf{m} = \mathbf{h}$ 

• Important property: for any NDF *D*, the "sum" of the area of the micro-facets, when projected on outgoing direction **v**, must equal the area of the macro-surface, projected on **v** 

$$(\mathbf{v} \cdot \mathbf{n}) = \int_{\Omega} D(\mathbf{m}_{\omega})(\mathbf{v} \cdot \mathbf{m}) d\omega$$



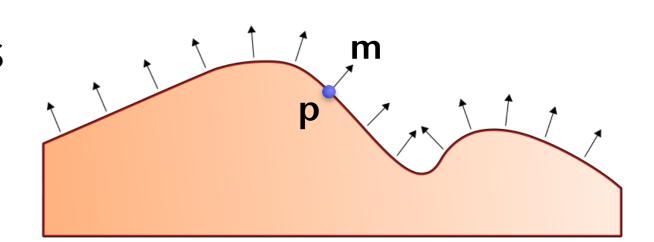
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## **Another Important Property**



• The NDF allows us to switch from spatial integrals over the surface to "statistical" integrals over the hemisphere of normals  $\Omega$ 



 The Gauß map: maps every point on a surface to its normal

$$G: \mathbf{p} \in M \mapsto \mathbf{m}_p \in \Omega$$



• Consider  $\Omega' \subseteq \Omega$ ; this induces

$$M' \subseteq M$$
 with  $M' = \{ \mathbf{p} \in M \mid \mathbf{m}(\mathbf{p}) \in \Omega' \}$ 

• The NDF *D* has (must have) the property (sketch of proof later):

$$\int_{M'} dp = \int_{\Omega'} D(\mathbf{m}) dm \tag{*}$$



#### Nice consequences



The area can be computed two ways:

Area
$$(M) = \int_{M} dp = \int_{O} D(\mathbf{m}) dm$$

 Integrals over the surface can be converted into integrals over the hemisphere:

$$\int_{M} f(\mathbf{m}(\mathbf{p})) d\mathbf{p} = \int_{\Omega} f(\mathbf{m}) D(\mathbf{m}) d\mathbf{m}$$

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## **FYI:** Sketch of Proof for Equation (\*)



- Let M be the micro-surface, p a point on M
- From differential geometry:

$$\kappa(\mathbf{p}) = \lim_{A \to 0} \frac{A'}{A}$$

where A = area of small patch around  $\mathbf{p} \in M$ , shrinking smaller and smaller,  $A' = G(A) = \{\mathbf{m}(\mathbf{p}) \mid \mathbf{p} \in A\} \subseteq \Omega$  (Gauß map of A) is a small area on  $\Omega$ 

- For small A, we can also write  $\kappa(\mathbf{p}) = \frac{dA'}{dA}$ , or  $dA = \frac{1}{\kappa(\mathbf{p})} dA'$
- Replacing dA' (area on  $\Omega$ ) with  $d\omega_p$  (solid angle), we can replace integrals over micro surfaces by integrals over  $\Omega$ :

$$\int_{M} f(\mathbf{m}(\mathbf{p})) dA = \int_{\Omega} f(\mathbf{m}) \frac{1}{\kappa(\mathbf{p})} d\omega$$

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#### Derivation of an NDF

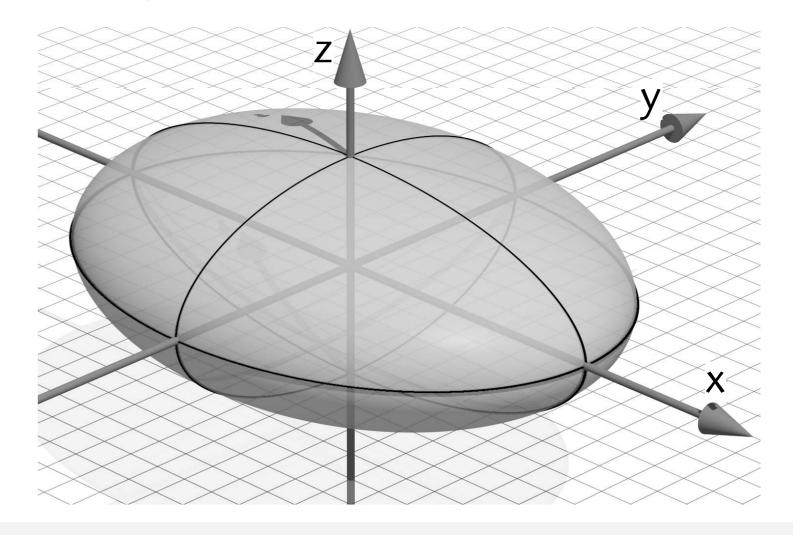


- Assumption: normals on M are distributed like on an ellipsoid
- General (implicit) equation of general ellipsoids:

$$f(\mathbf{p}) = \mathbf{p}^{\mathsf{T}} A^{\mathsf{T}} A \mathbf{p} - 1 = 0$$

with  $A = R \cdot S$ , where R = rotation matrix, S = scaling

- Simplification: choose  $A = \begin{pmatrix} \alpha_x & \alpha_y \\ & 1 \end{pmatrix}$ 
  - $\alpha_x$  and  $\alpha_y$  are kind of "roughness" params
- Note: the following is only a sketch of a derivation lots of steps/theorems in between, which are
  rather deep results from differential geometry,
  have been omitted (otherwise, much more time
  would be needed in class)







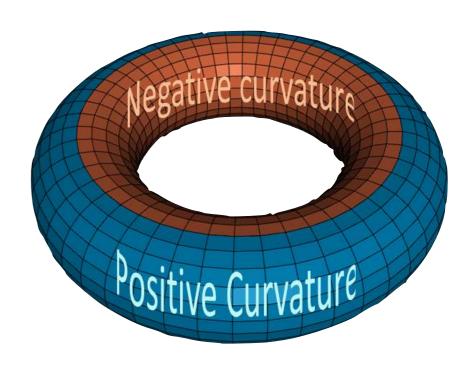
• From differential geometry: integrals over surfaces can be translated into integrals over the Gauss map, using the factor

$$D(\mathbf{m}) = \frac{1}{\dots \kappa(\mathbf{m})}$$

where  $\kappa$  = Gaussian curvature

(The "..." are normalization terms to maintain conservation of energy, omitted here for clarity)

Next goal: derive  $\kappa$  for our special surface (ellipsoid)







 More from differential geometry: implicit surfaces have the following Gaussian curvature at point p on the surface

$$\kappa(\mathbf{p}) = \frac{(\nabla f)^{\mathsf{T}} \mathrm{adj}(H) \nabla f}{||\nabla f||^4}$$
(1)

where

$$abla f$$
 = derivative of  $f=2A^2\mathbf{p}=H\mathbf{p}$ 
 $H=
abla^2f$  = Hessian matrix of  $f=2A^2$ 
adj $(H)$  = adjoint matrix of  $H=\det(H)H^{-1}=4\det(A)^2A^{-2}$ 





For all implicit surfaces, the normal m at a point p on the surface is

$$\mathbf{m} = \frac{\nabla f(\mathbf{p})}{||\nabla f||}$$

Plugging everything into (1):

$$\kappa(\mathbf{p}) = \frac{\det(H)\mathbf{m}^{\mathsf{T}}H^{-1}\mathbf{m}}{||\nabla f||^2}$$
(2)

- Remember: f is an ellipsoid, so there must be a 1:1 mapping between normals m (on the Gauss map) and points p (on the surface)
- Next goal: get rid of the  $||\nabla f||^2$  in the denominator





- Look at the term  $\mathbf{m}^{\mathsf{T}}H^{-1}\mathbf{m}$
- Use  $\mathbf{m} = \frac{\nabla f}{||\nabla f||}$ , plugging it into previous term, we get

$$\mathbf{m}^{\mathsf{T}}H^{-1}\mathbf{m} = \frac{(\nabla f)^{\mathsf{T}} \cdot H^{-1} \cdot \nabla f}{||\nabla f||^2} = \frac{(H\mathbf{p})^{\mathsf{T}} \cdot H^{-1} \cdot (H\mathbf{p})}{||\nabla f||^2} = \frac{\mathbf{p}^{\mathsf{T}}H^{\mathsf{T}}\mathbf{p}}{||\nabla f||^2}$$

$$= \frac{2\mathbf{p}^{\mathsf{T}}(A^2)^{\mathsf{T}}\mathbf{p}}{4||A^2\mathbf{p}||^2} = \frac{1}{2||A^2\mathbf{p}||^2}$$

Thus,

$$\frac{1}{||A^2\mathbf{p}||^2} = 2\mathbf{m}^\mathsf{T} H^{-1}\mathbf{m} = \mathbf{m}^\mathsf{T} A^{-2}\mathbf{m}$$

(3)

Physically-Based Lighting





• Now, we can rewrite (1):

$$\kappa(\mathbf{p}) = \frac{\det(H)\mathbf{m}^{\mathsf{T}}H^{-1}\mathbf{m}}{||\nabla f||^2} = \frac{2^3 \det(A)^2 \mathbf{m}^{\mathsf{T}} \frac{1}{2}A^{-2}\mathbf{m}}{2^2 ||A^2 \mathbf{p}||^2} = \det(A)^2 (\mathbf{m}^{\mathsf{T}}A^{-2}\mathbf{m})(\mathbf{m}^{\mathsf{T}}A^{-2}\mathbf{m})$$

- We can simplify this further by using  $\mathbf{m}^{\mathsf{T}} A^{-2} \mathbf{m} = (A^{-1} \mathbf{m})^2$ 
  - This is because A is special:  $m^{T}A^{-2}m = m^{T}(A^{-1})^{T}A^{-1}m = (A^{-1}m)^{T}A^{-1}m$
- In total:  $\kappa(m) = \det(A)^2 (A^{-1}m)^4$
- Thus:  $D(\mathbf{m}) = \frac{1}{\kappa(\mathbf{m})} = \frac{1}{(\alpha_x' \alpha_y')^2 (A'\mathbf{m})^4}$  with  $\alpha_x' = \frac{1}{\alpha_x}$ ,  $A' = A^{-1} = \text{diag}(\alpha_x', \alpha_y', 1)$





- Generalization: allow  $A = R \cdot S$ , with non-trivial rotation
  - Effect: predominant normals direction is *not* along macroscopic normal
  - Is it really necessary? would it be visible? are there such materials?
- Further specialization: choose  $\alpha_x = \alpha_y = \alpha$  (i.e., isotropic NDF) (note: here, I will write  $\alpha$ , instead of  $\alpha'$ , just for sake of simplicity)
  - Thus, we get

$$D(\mathbf{m}) = \frac{1}{\alpha^4 (A\mathbf{m})^4}$$
 with  $A = \text{diag}(\alpha, \alpha, 1)$ 

• You can use this equation as-is, or make it more obviously depend on  $\theta$ 





- Use polar coords to represent m:  $\mathbf{m} = \begin{pmatrix} \cos \varphi \sin \theta \\ \sin \varphi \sin \theta \\ \cos \theta \end{pmatrix}$
- Consider the term

$$(A\mathbf{m})^2 = \alpha^2(\cos\varphi\sin\theta)^2 + \alpha^2(\sin\varphi\sin\theta)^2 + (\cos\theta)^2$$
$$= \alpha^2(1 - \cos^2\theta) + \cos^2\theta$$

$$m \rightarrow \theta$$

(using twice 
$$cos^2 + sin^2 = 1$$
)

$$= \alpha^2 + (1 - \alpha^2)\cos^2\theta = \alpha^2(1 + (\frac{1}{\alpha^2} - 1)\cos^2\theta)$$
 (just a bit of terms rearrangement)

• Remember,  $\alpha$  is in reality  $\frac{1}{\alpha}$ , and  $\cos \theta = \mathbf{n} \cdot \mathbf{m}$ , so we stick both back in:

$$(Am)^2 = \frac{1}{\alpha^2} (1 + (\alpha^2 - 1)(nm)^2)$$

• In total, we get:  $D(\mathbf{m}) = \frac{1}{\frac{1}{\alpha^4} (A' \mathbf{m})^4} = \frac{\alpha^8}{\left(1 + (\alpha^2 - 1)(\mathbf{nm})^2\right)^2}$ 



# A Commonly Used NDF



Macrosurface

• Often, the so-called Trowbridge-Reitz/GGX is used (which is isotropic wrt.

roughness):

$$D(\mathbf{h}) = \frac{\alpha^2}{\pi((\mathbf{n} \cdot \mathbf{h})^2(\alpha^2 - 1) + 1)^2}$$

where  $\alpha$  = roughness

Often, artists prefer the following transfer function:

$$\alpha = (1 - \mathsf{glossiness})^2$$

 Glossiness is more intuitive (brighter texels in the "roughness" texture means "more shiny")

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Apparently, the square function is aesthetically more pleasing

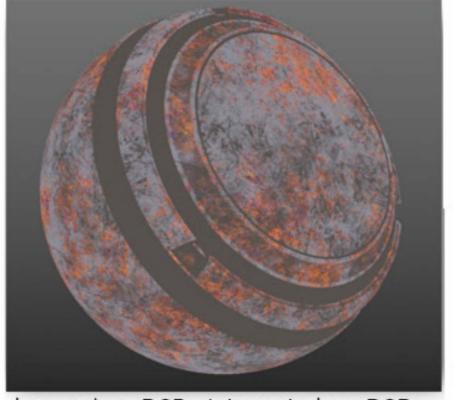
Microfacet



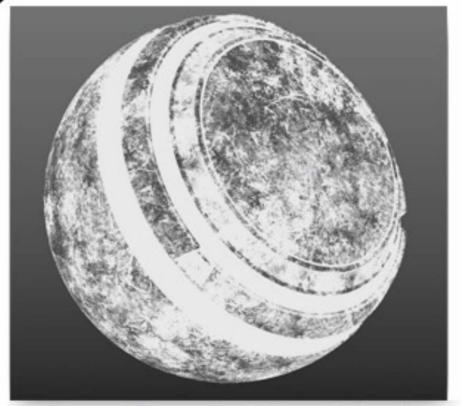
# Example of Roughness, Color, and Metallic Textures



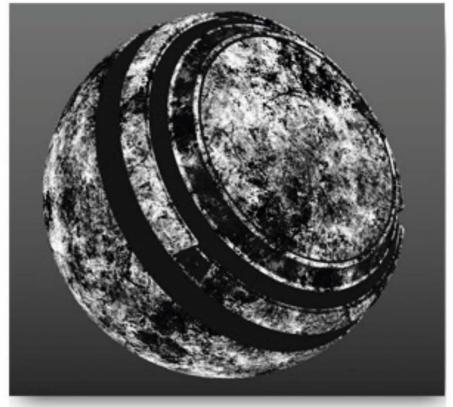




base color - RGB - interpreted as sRGB



roughness - Grayscale - interpreted as linear



metallic- Grayscale - interpreted as linear



# The Geometry Function



• The geometry function  $G = G(l, \mathbf{v}, \mathbf{h})$  gives the "likelihood" that micro-facets with  $\mathbf{m} = \mathbf{h} = \frac{1}{2}(l + \mathbf{v})$  are visible from the outgoing direction as well as for the incoming light



• G is symmetric in *l* and v

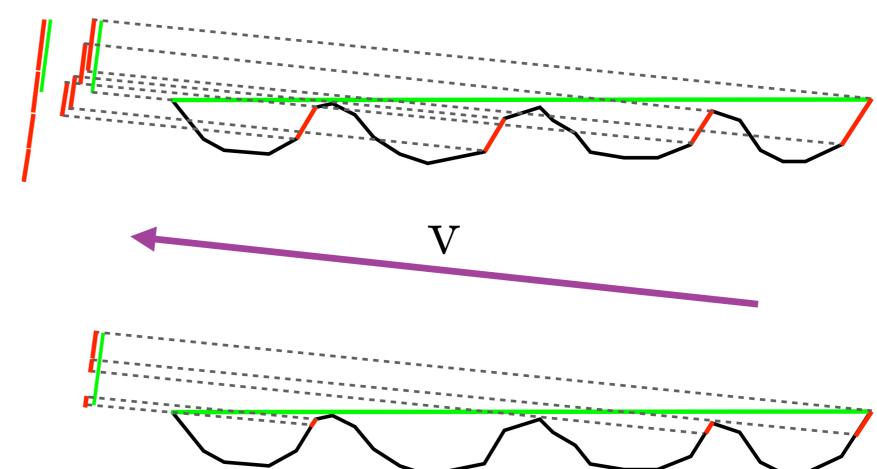


Commonly used function (Smith):

$$G(l, \mathbf{v}, \mathbf{h}) = G'(l)G'(\mathbf{v})$$

$$G'(\mathsf{x}) = rac{\mathsf{n} \cdot \mathsf{x}}{(\mathsf{n} \cdot \mathsf{x})(1-eta) + eta}$$

where  $\beta = \left(\frac{\text{roughness}+1}{2}\right)^2$  (again for aesthetic reasons)



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# A Current Standard is the "Disney BRDF"



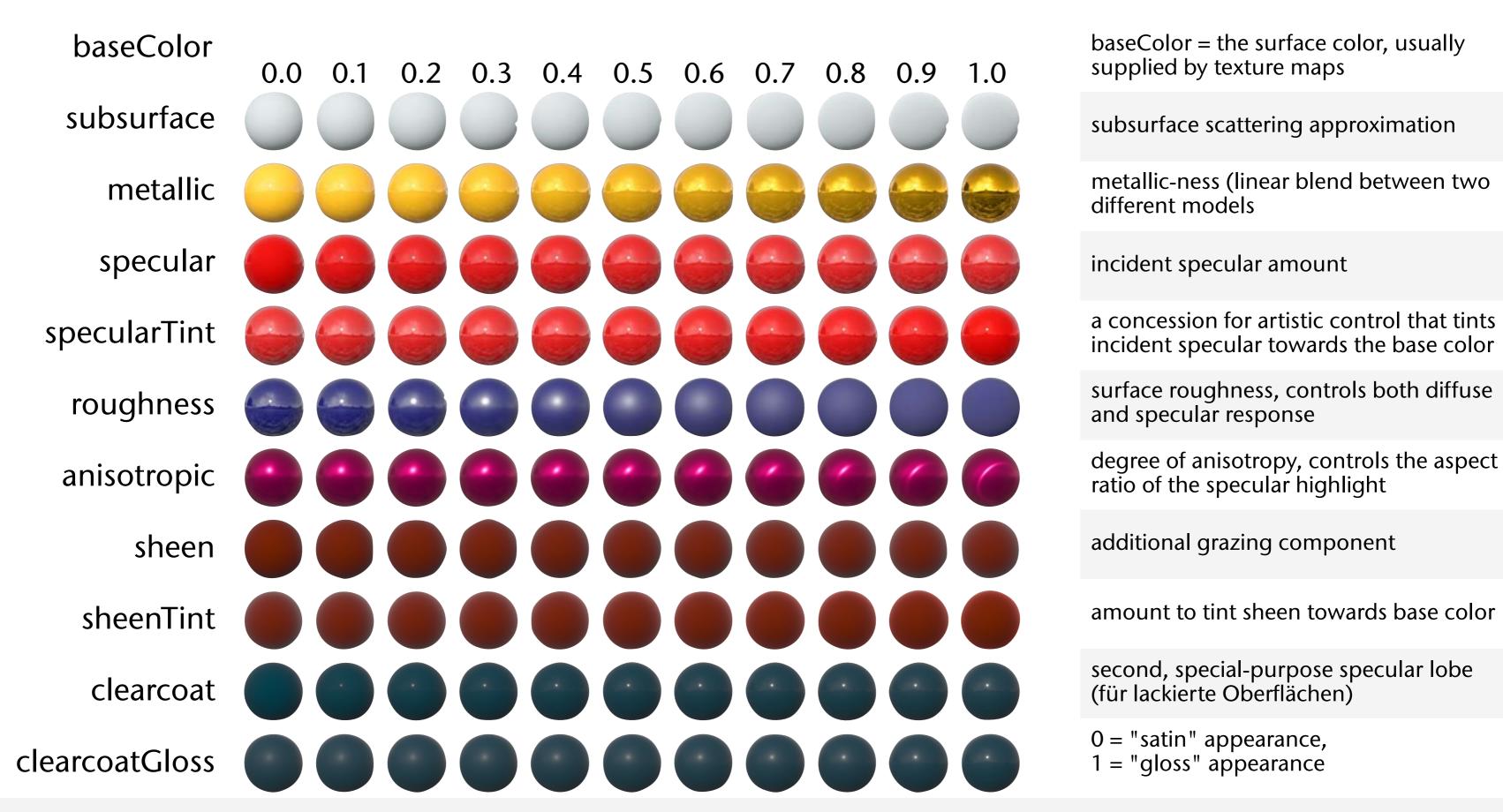
- Sometimes in general called metal/roughness workflow
- Design principles, in order to make work in the "art department" easier and more intuitive:
  - 1. Try to be physically correct, but intuitiveness of parameters has priority over correctness
  - 2. As few parameters as possible
  - 3. All parameter should be in [0,1]
  - 4. All parameters settings and combinations should be possible and lead to plausible results (no "funny" effects for specific values/combinations)
- Has still a lot more parameters than we have covered
  - But most can be modeled easily by introducing similar, additional terms



#### The Parameters



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# Sources / Literature



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- Naty Hoffman: Background: *Physics and Math of Shading*. 2013/2015
- Manfredo do Carmo: Differential Geometry. Prentice Hall, 1976.
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